

# A Hamiltonian System with an Even Term

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## 1. Introduction

In this paper we study, using variational methods, an equation of the form  $-u'' + u = h(t)V(u)$ , where  $h$  and  $V$  are differentiable,  $h$  is positive, bounded, and bounded away from zero, and  $V$  is a “superquadratic” potential. That is,  $V$  behaves like  $q$  to a power greater than 2, so  $|V(q)| = o(|q|^2)$  for  $|q|$  small and  $V(q) > O(|q|^2)$  for  $|q|$  large. To prove that a solution homoclinic to zero exists, one must assume additional hypotheses on  $h$  (see [EL] for a counterexample). In [R1], solutions were found when  $h$  is assumed to be periodic. In [STT], solutions were found when  $h$  is almost periodic (a weaker condition than periodicity). In [MNT], a condition yet weaker than almost periodic is defined, and solutions to the equation are found when  $h$  satisfies that condition. Like periodicity and almost periodicity, this condition assumes basically that  $h$  is similar to translates of itself, that is, for certain large values of  $T$ , the functions  $t \mapsto h(t)$  and  $t \mapsto h(t+T)$  are close to each other. Other ways to guarantee solutions involve making  $|h'|$  small: see papers such as [FW], [WZ], and [FdP] on the nonlinear Schrödinger equation, and [A] for a novel example of an  $h$  which “oscillates slowly.”

In this paper we attempt to find solutions to the equation without assuming that  $h$  satisfies any kind of time-recurrence property or restriction on  $h'$ . We assume two conditions: first, that  $h$  is even ( $h(-t) = h(t)$ ). Therefore it is convenient to treat the equation as an equation on the half-line  $\mathbf{R}^+ = [0, \infty)$ . Second,  $h$  only takes on a small range of values, with the variation in  $h$  depending on  $V$ . Here is a statement of the theorem:

**THEOREM 1.0** *Let  $n \geq 1$  and  $V$  satisfy*

(V<sub>1</sub>)  $V \in C^2(\mathbf{R}^n, \mathbf{R})$

(V<sub>2</sub>)  $V'(0) = 0$ ,  $V''(0) = 0$ , and

(V<sub>3</sub>) *there exists  $p > 1$  such that  $V''(q)q \cdot q \geq pV'(q)q > 0$  for all  $q \in \mathbf{R}^n \setminus \{0\}$ .*

*Then there exists  $d > 0$  with the property that if  $h$  satisfies*

(h<sub>1</sub>)  $h \in C^1(\mathbf{R}^+, \mathbf{R})$

(h<sub>2</sub>)  $h'(0) = 0$ , and

(h<sub>3</sub>)  $1 \leq h(t) \leq 1 + d$  for all  $t \in \mathbf{R}$ , then the differential equation

$$(*) \quad -u'' + u = h(t)V'(u)$$

has a non-zero solution  $v$  on  $\mathbf{R}^+$ , satisfying  $v'(0) = 0$  and  $v(t) \rightarrow 0$ ,  $v'(t) \rightarrow 0$  as  $t \rightarrow \infty$ .

An example of  $V$  satisfying (V<sub>1</sub>) – (V<sub>3</sub>) is  $V(q) = |q|^{p+1}$  with  $p > 1$ . Condition (V<sub>3</sub>) is a little stronger than growth conditions found in previous papers such as [Sé] or [CMN]. The conditions on  $h$  are fairly weak;  $h$  need not be periodic, or monotone, or tend to a single value as  $t \rightarrow \infty$  like in [BL]. If  $h$  has a lower bound other than 1, then  $h$  and  $V$  can be rescaled so that (h<sub>3</sub>) is satisfied and the problem reduces to the one in the theorem statement.

### Plan of Proof

We give a variational formulation of the problem. Let  $E = W^{1,2}(\mathbf{R}^+)$  along with the inner product

$$(u, w) = \int_0^\infty (u' \cdot w' + u \cdot w) dt$$

for  $u, w \in E$  and the associated norm  $\|u\| \equiv \|u\|_{W^{1,2}(\mathbf{R}^+)}$ . Then the functional  $I \in C^2(E, \mathbf{R})$  corresponding to (\*) is

$$I(u) = \frac{1}{2} \|u\|^2 - \int_0^\infty h(t)V(u(t)) dt.$$

Any critical point  $v$  of  $I$  satisfies the differential equation (\*), with  $v(t) \rightarrow 0$  and  $v'(t) \rightarrow 0$  as  $t \rightarrow \infty$ . Also, any critical point of  $I$  satisfies the boundary condition  $v'(0) = 0$ : suppose  $v$  is a critical point of  $I$ . Define  $h_2(t) = h(|t|)$  for  $t \in \mathbf{R}$ . Then, since  $h'(0) = 0$ ,  $h_2 \in C^1(\mathbf{R}, \mathbf{R})$ . Define the functional  $I_2$  on  $W^{1,2}(\mathbf{R})$  by  $I_2(u) = \frac{1}{2} \|u\|_{W^{1,2}(\mathbf{R})}^2 - \int_{\mathbf{R}} h_2(t)V(u(t)) dt$ , and  $v_2 \in W^{1,2}(\mathbf{R})$  by  $v_2(t) = v(|t|)$ . Then it is easy to verify that  $v_2$  is a critical point of  $I_2$ , and therefore a classical solution of the equation  $-u'' + u = h_2(t)V'(u)$  on the entire real line. Since  $h_2$  is an even function of  $t$ , and  $h_2 \in C^1(\mathbf{R})$ ,  $v_2'(0) = 0$ , so  $v'(0) = 0$ .

We will prove via an indirect argument that a critical point of  $I$  exists. First we define a submanifold  $\mathcal{S}$  of  $E = W^{1,2}(\mathbf{R}^+)$  with the property that  $\inf_{u \in \mathcal{S}} I(u) = c$ , where  $c$  is the mountain-pass value associated with  $I$ . Then we take a sequence  $(u_m) \subset E$  with  $I(u_m) \rightarrow c$  and  $I'(u_m) \rightarrow 0$  as  $m \rightarrow \infty$ . It is not apparent whether  $I$  satisfies the Palais-Smale condition, so it is not clear whether  $(u_m)$  converges. But we can show that  $(u_m)$  is a bounded sequence, so it has a weak limit. This weak limit point must be a critical point of  $I$ . If the limit point is not zero, then Theorem 1.0 is proven.

If  $(u_m)$  converges weakly to zero, then matters are more complicated. In this case, we can construct a sequence  $(y_m)$  with  $I(y_m) \leq c/2 + o(m)$ , where  $o(m) \rightarrow 0$  as  $m \rightarrow \infty$ , and  $y_m$  “close” to  $\mathcal{S}$ . For large enough  $m$ , we can use  $y_m$  to construct  $z \in \mathcal{S}$  with  $I(z) < c$ . This is impossible, so  $(u_m)$  has a nonzero weak limit, and there exists  $v$  satisfying Theorem 1.0.

### Organization of Paper

This paper is organized as follows: in Section 2 we explore the mountain-pass structure of the functional  $I$ , define the manifold  $\mathcal{S}$ , and obtain some quantitative estimates. Section 3 contains the main proof of Theorem 1.0, the “splitting” argument to obtain the sequence  $(y_m) \subset \mathcal{S}$  in the indirect argument above. Section 4 contains a computation of  $d$  for a specific function  $V$ .

## 2. Mountain-Pass Structure of $I$

Before defining  $\mathcal{S}$ , let us explore the related mountain-pass structure of  $I$ . Define the set of paths

$$(2.0) \quad \Gamma = \{ \gamma \in C([0, 1], E) \mid \gamma(0) = 0, I(\gamma(1)) < 0 \}.$$

Integrating  $(V_3)$  yields

$$(2.1) \quad V'(q) \cdot q \geq (p+1)V(q)$$

for all  $q \in \mathbf{R}^n$ . For  $\lambda > 1$ , the above implies

$$(2.2) \quad V(\lambda q) \geq \lambda^{p+1}V(q)$$

for all  $q \in \mathbf{R}^n$ . Thus it is easy to show that for any  $u \in E \setminus \{0\}$ ,  $I(\lambda u) \rightarrow -\infty$  as  $\lambda \rightarrow \infty$ , and  $\Gamma$  is well defined. Define the minimax value

$$(2.3) \quad c = \inf_{\gamma \in \Gamma} \max_{\theta \in [0, 1]} I(\gamma(\theta)).$$

Let us obtain a positive lower bound for  $c$ . Let  $\beta > 0$  satisfy

$$(2.4) \quad |q| \leq \beta \Rightarrow V'(q) \cdot q \leq \frac{1}{8}|q|^2.$$

This is possible by  $(V_1) - (V_2)$ . From now on assume, without loss of generality, that

$$(2.5) \quad d \leq 1.$$

Then  $h(t) \leq 2$  for all  $t \geq 0$ . If  $\|u\| \leq \beta$ , then  $\|u\|_{L^\infty(\mathbf{R}_+) } \leq \beta$  (see Appendix), and

$$(2.6) \quad \begin{aligned} I(u) &= \frac{1}{2}\|u\|^2 - \int_0^\infty h(t)V(u) dt \geq \frac{1}{2}\|u\|^2 - \frac{2}{p+1} \int_0^\infty V'(u) \cdot u dt \geq \\ &\text{(by (2.1))} \\ &\geq \frac{1}{2}\|u\|^2 - (1) \int_0^\infty \frac{1}{8}|u|^2 dt \geq \frac{1}{2}\|u\|^2 - \frac{1}{8}\|u\|^2 = \frac{3}{8}\|u\|^2 \geq 0. \end{aligned}$$

Therefore any mountain-pass curve must cross the sphere  $\{\|u\| = \beta\}$ , that is, if  $\gamma \in \Gamma$ , there exists  $\theta^* \in [0, 1]$  with  $\|\gamma(\theta^*)\| = \beta$ . So the above implies

$$(2.7) \quad \max_{\theta \in [0, 1]} I(\gamma(\theta)) \geq I(\gamma(\theta^*)) \geq \frac{3}{8}\|\gamma(\theta^*)\|^2 = \frac{3}{8}\beta^2.$$

Since  $\gamma$  is an arbitrary element of  $\Gamma$ ,

$$(2.8) \quad c \geq \frac{3}{8}\beta^2.$$

Note that this estimate does not depend on  $d$ , as long as  $d \leq 1$ .

There is another way to describe  $c$  (we will need both characterizations). Define

$$(2.9) \quad \mathcal{S} = \{u \in E \mid u \neq 0, I'(u)u = 0\}.$$

In [R2] it is proven, under weaker growth hypotheses on  $V$  than  $(V_3)$ , that

$$(2.10) \quad \inf_{u \in \mathcal{S}} I(u) = c.$$

In fact, for any  $u \in \mathcal{S}$ , the function  $s \mapsto I(su)$  is strictly increasing on  $0 < s < 1$ , attains a maximum of  $I(u)$  at  $s = 1$ , and decreases to  $-\infty$  on  $1 < s < \infty$ . The following lemma gives estimates how quickly  $I(su)$  changes when  $s$  is near 1:

LEMMA 2.11 *Let  $u \in E$  and define  $g(s) = I(su)$  for  $s \geq 0$ . Assume  $p \leq 2$ . Then*

$$(i) \quad s \geq 1 \Rightarrow g'(s) \leq g'(1)s^p - \frac{1}{4}(p-1)(s-1)\|u\|^2$$

and

$$(ii) \quad \frac{1}{2} \leq s \leq 1 \Rightarrow g'(s) \geq g'(1)s^p + \frac{1}{4}(p-1)(1-s)\|u\|^2.$$

*Proof:* Let  $u \in E$  and define  $g(s) = I(su)$ . Then

$$(2.12) \quad g(s) = \frac{1}{2}s^2\|u\|^2 - \int_0^\infty h(t)V(su) dt, \quad g'(s) = s\|u\|^2 - \int_0^\infty h(t)V'(su) \cdot u dt,$$

$$g''(s) = \|u\|^2 - \int_0^\infty h(t)V''(su)u \cdot u dt.$$

By  $(V_3)$ ,

$$(2.13) \quad g''(s) = \|u\|^2 - \frac{1}{s^2} \int_0^\infty h(t)V''(su)(su) \cdot (su) dt \leq \|u\|^2 - \frac{p}{s^2} \int_0^\infty h(t)V'(su) \cdot (su) dt =$$

$$= \|u\|^2 - \frac{p}{s} \int_0^\infty h(t)V'(su) \cdot u dt = \|u\|^2 - \frac{p}{s}(s\|u\|^2 - g'(s)) =$$

$$= \frac{p}{s}g'(s) - (p-1)\|u\|^2.$$

Therefore,

$$(2.14) \quad \frac{d}{ds}[s^{-p}g'(s)] = s^{-p}g''(s) - ps^{-p-1}g'(s) = s^{-p}(g''(s) - \frac{p}{s}g'(s)) \leq -(p-1)s^{-p}\|u\|^2.$$

If  $s \geq 1$ , then integrating the above from 1 to  $s$  yields

$$(2.15) \quad s^{-p}g'(s) - g'(1) \leq -(p-1)\|u\|^2 \int_1^s s^{-p} ds = -(1-s^{-p+1})\|u\|^2,$$

$$g'(s) \leq s^p g'(1) - (s^p - s)\|u\|^2.$$

If  $s \leq 1$ , then integrating (2.14) from  $s$  to 1 yields

$$(2.16) \quad g'(1) - s^{-p}g'(s) \leq -(p-1)\|u\|^2 \int_s^1 t^{-p} dt = (1-s^{-p+1})\|u\|^2,$$

$$g'(s) \geq s^p g'(1) + (s - s^p)\|u\|^2.$$

If  $s \geq 1$ , then by the mean value theorem, there exists  $\lambda \geq s \geq 1$  with

$$(2.17) \quad s^p - s \geq s^{p-1} - 1 \geq (p-1)\lambda^{p-2}(s-1) \geq (p-1)(t-1).$$

If  $s \in [1/2, 1]$ , then  $1/s \geq 1$ , so by the above,

$$(2.18) \quad \begin{aligned} s - s^p &= s^{p+1}(1/s^p - 1/s) \geq (p-1)s^{p+1}(1/s - 1) = (p-1)s^p(1-s) \\ &= \frac{1^p}{2}(p-1)(1-s) \geq \frac{1}{4}(p-1)(1-s). \end{aligned}$$

Lemma 2.11 follows from (2.15)-(2.18).

We have a lower bound for  $c$  that is independent of  $d$ . We also need an upper bound for  $c$  that is independent of  $d$ . Define the functional

$$(2.19) \quad I^+(u) = \frac{1}{2}\|u\|^2 - \int_0^\infty F(u(t)) dt.$$

Then  $I^+(u) \geq I(u)$  for all  $u \in E$ . Define the mountain-pass value  $c^+$ , similar to  $c$ , by defining the set of paths

$$(2.20) \quad \Gamma^+ = \{g \in C([0, 1], E) \mid g(0) = 0, I^+(g(1)) < 0\},$$

and setting

$$(2.21) \quad c^+ = \inf_{g \in \Gamma^+} \max_{\theta \in [0, 1]} I^+(g(\theta)).$$

$c^+$  depends only on  $V$ , not on  $d$ . Using the mountain-pass characterization of  $c$  ((2.3)), it is easy to see that  $c^+ \geq c$  because  $I^+(u) \geq I(u)$  for all  $u \in E$ . We will estimate  $c^+$  in terms of  $\beta$  and  $V$  in Section 4.

It is well known that  $(V_3)$  or a weaker condition implies that Palais-Smale sequences of  $I$  are bounded, even that  $\mathcal{S} \cap \{u \mid I(u) \leq D\}$  is bounded for any  $D \in \mathbf{R}$ . We want an estimate on  $\|u\|$  for when  $I(u)$  is small and  $u$  is ‘‘almost’’ in  $\mathcal{S}$ :

LEMMA 2.22 *If  $p \leq 2$ ,  $|I'(u)u| \leq c^+$  and  $I(u) \leq 2c^+$ , then*

$$(2.23) \quad \|u\| \leq \sqrt{\frac{14c^+}{p-1}} \equiv B.$$

Proof:

$$\begin{aligned} -c^+ &\leq I'(u)u = \|u\|^2 - \int_{\mathbf{R}} hV'(u) \cdot u \leq \|u\|^2 - (p+1) \int_{\mathbf{R}} hV(u) = \\ &= (p+1)I(u) - \left(\frac{p-1}{2}\right)\|u\|^2 \leq 6c^+ - \left(\frac{p-1}{2}\right)\|u\|^2 \end{aligned}$$

by (2.1), so

$$\|u\|^2 \leq \left(\frac{2}{p-1}\right) \cdot 7c^+ = \frac{14c^+}{p-1}.$$

### 3. Splitting

This section contains the ‘‘splitting’’ argument that is the core of the proof of Theorem 1.0. By Ekeland’s Variational Principle ([MW]), there exists a Palais-Smale sequence  $(u_m) \subset E$  with  $I(u_m) \rightarrow c$  and  $I'(u_m) \rightarrow 0$

as  $m \rightarrow \infty$ . By arguments of [CR],  $(u_m)$  is bounded. Therefore it has a subsequential weak limit  $\bar{u}$ . Also by [CR],  $\bar{u}$  is a critical point of  $I$ , and  $u_m$  converges to  $\bar{u}$  in  $W^{1,2}([0, R])$  for each  $R > 0$ . If  $\bar{u} \neq 0$ , then Theorem 1.0 is proven. In fact, in this case,  $I(\bar{u}) \leq c$  (see [CR]).  $I(\bar{u}) \geq c$  because by the observations following (2.10), for large enough  $T$ ,  $\theta \mapsto T\theta\bar{u}$  defines a path in  $\Gamma$ , along which the maximum value of  $I$  is  $c$ . Thus  $I(\bar{u}) = c$ .

We will show that if  $d$  is chosen small enough, in terms of  $V$ , then the case  $\bar{u} = 0$  is impossible. The argument is indirect. Suppose  $\bar{u} = 0$ . Define the cutoff function  $\varphi \in C(\mathbf{R}^+, [0, 1])$  by  $\varphi(t) = t$  for  $0 \leq t \leq 1$ ,  $\varphi \equiv 1$  on  $[1, \infty)$ . Define  $w_m = \varphi u_m$ .  $\|u_m\|_{W^{1,2}([0,1])} \rightarrow 0$  as  $m \rightarrow \infty$ , and it is easy to verify that  $\|u_m - w_m\| \rightarrow 0$  as  $m \rightarrow \infty$ .  $I''$ ,  $I'$ , and  $I$  are bounded on bounded subsets of  $E$ . For example, to prove for  $I''$ , let  $K > 0$  and suppose  $\|u\| \leq K$ . Then  $\|u\|_{L^\infty(\mathbf{R}^+)} \leq K$  (see Appendix). Let  $C > 0$  satisfy  $|V''(q)x \cdot y| \leq C$  for all  $|q| \leq K$ ,  $|x| \leq 1$ ,  $|y| \leq 1$ . Let  $v, w \in E$ . Then

$$(3.0) \quad \begin{aligned} |I''(u)(v, w)| &= |(v, w) - \int_0^\infty h(t)V''(u)v \cdot w dt| \leq \|v\|\|w\| + \int_0^\infty 2C|v||w| dt \leq \\ &\leq \|v\|\|w\| + 2C\|v\|_{L^2(\mathbf{R}^+)}\|w\|_{L^2(\mathbf{R}^+)} \leq (1 + 2C)\|v\|\|w\|. \end{aligned}$$

Since  $I''$ ,  $I'$ , and  $I$  are bounded on bounded subsets of  $E$ , and  $(u_m)$  is a bounded sequence, it follows that  $I(w_m) \rightarrow c$  and  $I'(w_m)w_m \rightarrow 0$  as  $m \rightarrow \infty$ .

Let  $\epsilon > 0$  satisfy

$$(3.1) \quad \epsilon < \beta^2/4$$

where  $\beta$  is from (2.4).  $\epsilon$  will fixed more precisely later. Since  $w_m \rightarrow 0$  in  $W^{1,2}([0, 1])$  (and thus in  $L^\infty([0, 1])$ ), we may choose  $m$  large enough so that

$$(3.2) \quad \|w_m\|_{L^\infty([0,1])} < \beta,$$

$$(3.3) \quad I(w_m) < \frac{7}{6}c,$$

and

$$(3.4) \quad |I'(w_m)w_m| < \epsilon.$$

For convenience define

$$(3.5) \quad w = w_m.$$

We will choose a ‘‘cutting point’’  $\hat{t} > 0$ , and split  $w$  into two functions,  $w^{(1)} = w|_{[0, \hat{t}]}$  (the restriction of  $w$  to  $[0, \hat{t}]$ ), and  $w^{(2)} = w|_{[\hat{t}, \infty)}$ .  $w^{(1)}$  and  $w^{(2)}$  can be transformed into functions  $z_1$  and  $z_2$  respectively in  $E$ :  $w^{(1)}$  into  $z_1$ , by reflecting over  $t = \hat{t}/2$ ; and  $w^{(2)}$  into  $z_2$ , by translating by a factor of  $\hat{t}$  to the left. If  $d$  is small enough and  $\hat{t}$  is chosen carefully,  $I'(z_1)z_1$  and  $I'(z_2)z_2$  are both very close to zero, but either  $I(z_1)$  or  $I(z_2)$  is significantly less than  $c$ . Using Lemma 2.11, we then choose  $\bar{s}$  very close to 1 so that  $\bar{s}z_* \in \mathcal{S}$  but  $I(\bar{s}z_*) < c$ , where  $*$  = 1 or 2. This contradicts the fact that  $\inf\{I(u) \mid u \in \mathcal{S}\} = c$ , proving Theorem 1.0.

Let us choose  $\hat{t}$ . We claim that  $\|w_m\|_{L^\infty(\mathbf{R}^+)} > \beta$  for large  $m$ : since  $I(w_m) \rightarrow c$  and  $I(0) \neq c$ ,  $\|w_m\|$  is bounded away from 0 for large  $m$ . If  $\|w_m\|_{L^\infty(\mathbf{R}^+)} \leq \beta$ , then by (2.4),

$$(3.6) \quad I'(w_m)(w_m) = \|w_m\|^2 - \int_0^\infty hV'(w_m) \cdot w_m \geq \|w_m\|^2 - \int_0^\infty 2\left(\frac{1}{8}\right)|w_m|^2 \geq \frac{3}{4}\|w_m\|^2.$$

This cannot happen for large  $m$ , since  $\|w_m\|$  is bounded away from 0 for large  $m$  and  $I'(w_m)w_m \rightarrow 0$ . Since  $\|w_m\|_{L^\infty(\mathbf{R}^+)} > \beta$  for large  $m$ , we may define

$$(3.7) \quad t_0 = \min\{t \mid |w(t)| \geq \beta\} < t_1 = \max\{t \mid |w(t)| \geq \beta\}.$$

By (3.2),  $1 < t_0 < t_1$ . By (3.4),

$$(3.8) \quad |I'(w)w| = \left| \int_0^\infty |w'|^2 + |w|^2 - h(t)V'(w) \cdot w \, dt \right| < \epsilon.$$

We will choose the cutting point  $\hat{t}$  between  $t_0$  and  $t_1$  so that the integral above, evaluated only from 0 to  $\hat{t}$ , is zero (and the integral evaluated from  $\hat{t}$  to  $\infty$  is also close to zero). For  $t < t_0$ ,  $|w(t)| < \beta$ , and since  $h(t) \leq 2$  for all  $t \geq 0$  ((2.5)),

$$(3.9) \quad |h(t)V'(w(t)) \cdot w(t)| \leq 2\left(\frac{1}{8}|w(t)|^2\right) = \frac{1}{4}|w(t)|^2$$

by the definition of  $\beta$  ((2.4)). Therefore

$$(3.10) \quad \begin{aligned} \int_0^{t_0} |w'|^2 + |w|^2 - h(t)V'(w(t)) \cdot w(t) \, dt &\geq \frac{3}{4} \int_0^{t_0} |w'|^2 + |w|^2 \, dt = \frac{3}{4} \|w\|_{W^{1,2}([0,t_0])}^2 \geq \\ &\geq \frac{3}{4} \|w\|_{W^{1,2}([0,t_0])}^2 \geq \frac{3}{16} \|w\|_{L^\infty([0,t_0])}^2 = \frac{3}{16} \beta^2, \end{aligned}$$

using an embedding in the Appendix, and the fact that  $\|w\|_{L^\infty([0,t_0])} = \beta$ . By similar reasoning to (3.9)-(3.10), and using the other embedding in the Appendix,

$$(3.11) \quad \int_{t_1}^\infty |w'|^2 + |w|^2 - h(t)V'(w(t)) \cdot w(t) \, dt \geq \frac{3}{4} \|w\|_{W^{1,2}([t_1,\infty])}^2 \geq \frac{3}{4} \|w\|_{L^\infty([t_1,\infty])}^2 = \frac{3}{4} \beta^2.$$

By (3.8), (3.11), and (3.1),

$$(3.12) \quad \begin{aligned} \int_0^{t_1} |w'|^2 + |w|^2 - h(t)V'(w(t)) \cdot w(t) \, dt &= \\ &= \int_0^\infty |w'|^2 + |w|^2 - h(t)V'(w) \cdot w \, dt - \int_{t_1}^\infty |w'|^2 + |w|^2 - h(t)V'(w) \cdot w \, dt < \\ &< \epsilon - \frac{3}{4} \beta^2 < \beta^2/4 - \frac{3}{4} \beta^2 < 0. \end{aligned}$$

The above integral is negative but the integral from 0 to  $t_0$  of the same integrand is positive ((3.10)). Therefore there exists  $\hat{t} \in (t_0, t_1)$  with

$$(3.13)(i) \quad \int_0^{\hat{t}} |w'|^2 + |w|^2 - h(t)V'(w(t)) \cdot w(t) \, dt = 0.$$

By the above and (3.8), we have, similarly,

$$(3.13)(ii) \quad \left| \int_{\hat{t}}^\infty |w'|^2 + |w|^2 - h(t)V'(w(t)) \cdot w(t) \, dt \right| < \epsilon.$$

By (3.3),

$$(3.14)(i) \quad \int_0^{\hat{t}} \frac{1}{2}|w'|^2 + \frac{1}{2}|w|^2 - h(t)V'(w(t)) \cdot w(t) \, dt < \frac{7}{12}c$$

or

$$(3.14)(ii) \quad \int_{\hat{t}}^{\infty} \frac{1}{2}\dot{w}^2 + \frac{1}{2}w^2 - h(t)F(w(t)) dt < \frac{7}{12}c.$$

If the former case, (3.14)(i), holds, define  $z \in E$  by reflecting  $w$  over  $t = \hat{t}/2$ , that is,

$$(3.15) \quad z(t) = \begin{cases} w(\hat{t} - t); & 0 \leq t \leq \hat{t} \\ 0; & t \geq \hat{t}. \end{cases}$$

If the latter case, (3.14)(ii), holds, define  $z \in E$  by  $z(t) = w(t + \hat{t})$ . In future arguments, we assume for convenience that the latter case holds. Arguments for the former case are very similar.

By the discussion preceding Lemma 2.11, there exists a unique  $\bar{s} > 0$  with the property that  $\bar{s}z \in \mathcal{S}$ . We will prove that, if one assumes  $d$  to be small enough, then  $I(\bar{s}z) < c$ . This is impossible, and Theorem 1.0 follows. Recall  $\epsilon$  from (3.1), and define  $\epsilon$  more precisely by

$$(3.16) \quad \epsilon = \frac{(p-1)\beta^2}{60}.$$

Set

$$(3.17) \quad d = \frac{\epsilon}{B^2} = \frac{(p-1)\beta^2}{60} \cdot \frac{(p-1)}{14c^+} = \frac{(p-1)^2\beta^2}{840c^+}.$$

Assume from now on that

$$(3.18) \quad p \leq 2.$$

Then, as we have been assuming,  $d \leq 1$ , using (2.8) and  $c^+ \geq c$ . The following estimate, which uses (2.8), will be useful later:

$$(3.19) \quad \epsilon = \frac{(p-1)\beta^2}{60} \leq \frac{(p-1)}{60} \cdot \frac{8}{3}c < \frac{(p-1)c}{22} \leq \frac{c}{22} \leq \frac{c^+}{22}.$$

We will show that  $|I'(z)z| < 3\epsilon$ , while  $I(z) < \frac{2}{3}c$ . This will imply that the function  $g(s) = I(sz)$  has a maximum for  $s \geq 0$  that is less than  $c$ , which is impossible. We estimate  $I'(z)z$  by comparing the integral for  $I'(z)z$  to that for  $I'(w)w$  in (3.13)(ii):

$$(3.20) \quad \begin{aligned} |I'(z)z| &= \left| \int_0^{\infty} |z'(t)|^2 + |z(t)|^2 - h(t)V'(z(t)) \cdot z(t) dt \right| = \\ &= \left| \int_0^{\infty} |w'(t + \hat{t})|^2 + |w(t + \hat{t})|^2 - h(t)V'(w(t + \hat{t})) \cdot w(t + \hat{t}) dt \right| = \\ &= \left| \int_{\hat{t}}^{\infty} |w'(t)|^2 + |w(t)|^2 - h(t - \hat{t})V'(w(t)) \cdot w(t) dt \right| \leq \\ &= \left| \int_{\hat{t}}^{\infty} |w'(t)|^2 + |w(t)|^2 - h(t)V'(w(t)) \cdot w(t) dt \right| + \\ &\quad + \left| \int_{\hat{t}}^{\infty} (h(t) - h(t - \hat{t}))V'(w(t)) \cdot w(t) dt \right| \leq \\ &= \epsilon + d \int_{\hat{t}}^{\infty} V'(w(t)) \cdot w(t) \leq \epsilon + d \int_0^{\infty} V'(w(t)) \cdot w(t) = \epsilon + d(\|w\|^2 - I'(w)w) \leq \\ &\quad \text{(by (3.13)(ii) and (V}_3\text{))} \\ &\leq \epsilon + d(B^2 + \epsilon) \leq 2\epsilon + dB^2 \leq 3\epsilon. \end{aligned}$$



In the last line we use (2.5) ( $d \leq 1$ ), and Lemma 2.22 with (3.3), (3.4) and (3.19).

Now we estimate  $I(z)$  by comparing the integral for  $I(z)$  to that for  $I(w)$ ; we assume case (3.14)(ii) holds, so  $z$  equals  $w$  translated  $\hat{t}$  units to the left. Recall that  $w$  satisfies (3.2)-(3.4).

$$\begin{aligned}
(3.21) \quad I(z) &= \int_0^\infty \frac{1}{2}|z'(t)|^2 + \frac{1}{2}|z(t)|^2 - h(t)V(z(t)) dt = \\
&= \int_0^\infty \frac{1}{2}|w'(t+\hat{t})|^2 + \frac{1}{2}|w(t+\hat{t})|^2 - h(t)V(w(t+\hat{t})) dt = \\
&= \int_{\hat{t}}^\infty \frac{1}{2}|w'(t)|^2 + \frac{1}{2}|w(t)|^2 - h(t-\hat{t})V(w(t)) dt = \\
&= \int_{\hat{t}}^\infty \frac{1}{2}|w'(t)|^2 + \frac{1}{2}|w(t)|^2 - h(t)V(w(t)) dt + \int_{\hat{t}}^\infty (h(t) - h(t-\hat{t}))V(w(t)) dt < \\
&< \frac{7}{12}c + d \int_{\hat{t}}^\infty V(w(t)) dt \leq \frac{7}{12}c + d \int_{\hat{t}}^\infty V'(w(t)) \cdot w(t) dt \leq \\
&\text{(by (3.3) and (2.1))} \\
&\leq \frac{7}{12}c + d(B^2 + \epsilon) < \frac{7}{12}c + 2\epsilon < \frac{2}{3}c.
\end{aligned}$$

In the last line, we estimate the last integral using the calculation at the end of (3.20), and also use (3.16),  $d \leq 1$ , and (3.19).

We have  $z \in E$  with  $I(z) < \frac{2}{3}c$  and  $|I'(z)z| < 3\epsilon$ . By choice of the cutting point  $\hat{t}$  between  $t_0$  and  $t_1$  ((3.7)), and the definition of  $z$  as a reflection or translation of  $w$  (see (3.15) and the remark following it),  $\|z\|_{L^\infty(\mathbf{R}^+)} \geq |z(0)| = \beta$ , so  $\|z\| \geq \beta$ . Defining  $g(s) = I(sz)$  as in Lemma 2.11,  $g(1) = I(z) < \frac{2}{3}c$  and  $|g'(1)| = |I'(z)z| < 3\epsilon$ . We will show that  $g'(5/4) < 0$  and  $g'(3/4) > 0$ . Therefore there exists  $\bar{s} \in (3/4, 5/4)$  with  $g'(\bar{s}) = I'(\bar{s}z)z = 0$ , so  $\bar{s}z \in \mathcal{S}$ . Then we prove that for all  $s \in [3/4, 5/4]$ ,  $g(s) < c$ . This contradicts the fact that  $I(\bar{s}z) \geq c$ , proving Theorem 1.0. By Lemma 2.11(i), since  $p \in (1, 2]$  and  $\|z\| > \beta$ ,

$$(3.22) \quad g'(5/4) \leq g'(1) - \frac{1}{4}(p-1)\left(\frac{1}{4}\right)\|z\|^2 \leq 3\epsilon - \frac{1}{16}(p-1)\beta^2 < 0$$

using the definition of  $\epsilon$  ((3.16)). Similarly,

$$(3.23) \quad g'(3/4) \geq g'(1) + \frac{1}{4}(p-1)\left(\frac{1}{4}\right)\|z\|^2 \geq -3\epsilon + \frac{1}{16}(p-1)\beta^2 > 0.$$

$|g'(1)| < 3\epsilon$ , so for  $s \in [1, 5/4]$ , Lemma 2.11(i) gives

$$(3.24) \quad g'(s) \leq g'(1)s^p - \frac{1}{2}(p-1)(s-1)\|z\|^2 \leq g'(1)s^p < 3\epsilon s^p < 3\epsilon(5/4)^2 < 5\epsilon,$$

and

$$(3.25) \quad g(s) = g(1) + \int_1^s g'(r) dr < \frac{2}{3}c + 5\epsilon(s-1) < \frac{2}{3}c + 2\epsilon < \frac{2}{3}c + \frac{c}{11} < c$$

(see (3.19)). For  $s \in [3/4, 1]$ , Lemma 2.11(ii) gives,

$$(3.26) \quad g'(s) \geq g'(1)s^p + \frac{1}{4}(p-1)(s-1)\|z\|^2 \geq g'(1)s^p > -3\epsilon s^p > -3\epsilon(1)^2 = -3\epsilon,$$

so

$$(3.27) \quad g(s) = g(1) - \int_s^1 g'(r) dr < \frac{2}{3}c + 3\epsilon(1-s) < \frac{2}{3}c + \epsilon < \frac{2}{3}c + c/22 < c$$

by (3.19). Therefore  $g(s) = I(sz) < c$  for all  $s \in [3/4, 5/4]$ . This is impossible because  $\bar{s}z \in \mathcal{S}$  for some  $\bar{s} \in [3/4, 5/4]$ . The assumption made at the beginning of this section is false. Theorem 1.0 is proven.

#### 4. Determining $d$ : an example.

Here we find how to write  $d$ , satisfying Theorem 1.0, compactly in terms of  $\beta$ ,  $p$ , and  $V$ . Then we find  $d$  for a specific function  $V$ .

To compute  $d$  using (3.17) we must estimate  $c^+$  as defined in (2.21). Let us find a way to estimate  $c^+$  for any  $V$  satisfying  $(V_1) - (V_3)$  and write it compactly. Recall  $I^+$ ,  $\Gamma^+$ , and  $c^+$  from (2.19)-(2.21). To define  $c^+$ , it suffices to find one element  $\gamma$  of  $\Gamma^+$  and choose  $c^+$  large enough to guarantee that  $c^+ \geq \max_{\theta > 0} I^+(g(\theta))$ . Define  $\beta$  as in (2.4). Let  $\bar{e}_1$  denote the unit vector  $[1 \ 0 \ 0 \ \cdots \ 0]^T \in \mathbf{R}^n$ , and define  $w : \mathbf{R}^+ \rightarrow \mathbf{R}$  by

$$(4.0) \quad w(t) = \beta e^{-t} \bar{e}_1.$$

A direct calculation yields  $\|w\| = \beta$ . Since  $\|w\|_{L^\infty(\mathbf{R}^+)} = \beta$ ,  $I^+(sw)(w) > 0$  for all  $s \in (0, 1]$ , by (3.6). Thus  $I(sw) < I(w)$  for all  $s \in (0, 1)$ . By (2.2),

$$(4.1) \quad I^+(sw) = \frac{1}{2} s^2 \|w\|^2 - \int_0^\infty V(sw) dt \leq \frac{1}{2} s^2 \beta^2 - s^{p+1} \int_0^\infty V(w) dt$$

for all  $s > 1$ .  $V(r\bar{e}_1)$  is increasing for positive  $r$ , so

$$(4.2) \quad \begin{aligned} \int_0^\infty V(w) dt &\geq \int_0^{\ln 2} V(w) dt = \int_0^{\ln 2} V(\beta e^{-s} \bar{e}_1) ds > \\ &> \int_0^{\ln 2} V(\beta \bar{e}_1 / 2) dt = (\ln 2) V(\beta \bar{e}_1 / 2) > V(\beta \bar{e}_1 / 2) / 2. \end{aligned}$$

Therefore

$$(4.3) \quad I^+(sw) \leq \alpha(s) \equiv \frac{1}{2} s^2 [\beta^2 - (F(\beta/2)/2) s^{p-1}]$$

for  $s > 1$ . By elementary calculus,  $\alpha(s)$  achieves a maximum over  $\{s > 0\}$  of

$$(4.4) \quad \frac{\beta^2}{2} \left( \frac{p-1}{p+1} \right) \left( \frac{4\beta^2}{(p+1)V(\beta/2)} \right)^{\frac{2}{p-1}} \leq \frac{\beta^2}{6} \left( \frac{2\beta^2}{V(\beta/2)} \right)^{\frac{2}{p-1}}.$$

The last expression is an upper bound for  $c^+$ . Using (3.17),  $d$  can be estimated by

$$(4.5) \quad \frac{(p-1)^2 \beta^2}{840 c^+} \geq \frac{(p-1)^2 \beta^2}{840} \cdot \frac{6}{\beta^2} \cdot \left( \frac{V(\beta \bar{e}_1 / 2)}{2\beta^2} \right)^{\frac{2}{p-1}} = \frac{(p-1)^2}{140} \left( \frac{V(\beta \bar{e}_1 / 2)}{2\beta^2} \right)^{\frac{2}{p-1}} \geq d.$$

Let us compute  $d$  for the specific case  $n = 1$ ,  $1 < p \leq 2$ ,  $V(q) = \frac{1}{p+1} |q|^{p+1}$ . We can pick  $\beta = \frac{1}{8} \frac{1}{p-1}$ , because

$$(4.6) \quad V'(q)q = |q|^{p+1} = |q|^{p-1} |q|^2 \leq \beta |q|^2$$

for  $|q| \leq \beta$ . Now,

$$(4.7) \quad V(\beta/2) = \frac{1}{p+1} \left( \frac{1}{8} \right)^{\frac{p+1}{p-1}} \geq \frac{1}{3} \left( \frac{1}{8} \right)^{\frac{3}{p-1}},$$

so, using (4.5),  $d$  can be estimated by

$$\begin{aligned} \frac{(p-1)^2}{140} \left( \frac{V(\beta/2)}{2\beta^2} \right)^{\frac{2}{p-1}} &\geq \frac{(p-1)^2}{140} \left( \frac{1}{6 \cdot 8^{\frac{3}{p-1}} \cdot 8^{\frac{2}{p-1}}} \right)^{\frac{2}{p-1}} > \\ &> \frac{(p-1)^2}{140} \left( \frac{1}{8} \right)^{\frac{2(p+4)}{(p-1)(p-1)}} \geq \frac{(p-1)^2}{140} \left( \frac{1}{8} \right)^{\frac{12}{(p-1)^2}} \geq d. \end{aligned}$$

## Appendix

This brief appendix contains two well-known Sobolev inequalities, along with embedding constants.

LEMMA 1 *If  $u \in W^{1,2}([0, \infty))$  then*

$$(i) \quad \|u\|_{L^\infty([0, \infty))} \leq \|u\|_{W^{1,2}([0, \infty))}.$$

*If  $a \geq 1$  and  $u \in W^{1,2}([0, a])$ , then*

$$(ii) \quad \|u\|_{L^\infty([0, a])} \leq 2\|u\|_{W^{1,2}([0, \infty))}.$$

*Proof of (i):* let  $u \in W^{1,2}([0, \infty))$  and  $x_1 \in [0, \infty)$ . Let  $\epsilon > 0$ . Choose  $x_0 \in [0, \infty)$  with  $|u(x_0)| < \epsilon$ . Then

$$\begin{aligned} u(x_1)^2 &= u(x_0)^2 + (u(x_1)^2 - u(x_0)^2) < \epsilon^2 + \left| \int_{x_0}^{x_1} \frac{d}{dx} u^2 dx \right| = \\ &= \epsilon^2 + \left| \int_{x_0}^{x_1} 2uu' dx \right| \leq \epsilon^2 + \left| \int_{x_0}^{x_1} (u')^2 + u^2 dx \right| \leq \epsilon^2 + \|u\|_{W^{1,2}([0, \infty))}^2 \end{aligned}$$

via the Cauchy-Schwarz inequality. Letting  $\epsilon$  go to zero,  $|u(x_1)| \leq \|u\|_{W^{1,2}([0, \infty))}$ . Since  $x_1$  is arbitrary, (i) is proven.

*Proof of (ii):* Let  $a \geq 1$  and  $u \in W^{1,2}([0, a])$ . Assume  $\|u\|_{L^\infty([0, a])} \geq 1$ . We will show that  $\|u\|_{W^{1,2}([0, \infty))} \geq 1/2$ .

If  $|u(x)| > 1/2$  for all  $x \in [0, a]$ , then  $\|u\|_{W^{1,2}([0, \infty))}^2 \geq \int_0^a u^2 > a/4 \geq 1/4$ . So suppose  $|u(x_0)| \leq 1/2$  for some  $x_0 \in [0, a]$ . Let  $x_1 \in [0, a]$  with  $|u(x_1)| \geq 1$ . Arguing as in part (i) above,

$$\begin{aligned} 1 \leq u(x_1)^2 &= u(x_0)^2 + (u(x_1)^2 - u(x_0)^2) < \frac{1}{4} + \left| \int_{x_0}^{x_1} \frac{d}{dx} u^2 dx \right| = \\ &= \frac{1}{4} + \left| \int_{x_0}^{x_1} 2uu' dx \right| \leq \frac{1}{4} + \left| \int_{x_0}^{x_1} (u')^2 + u^2 dx \right| \leq \frac{1}{4} + \|u\|_{W^{1,2}([0, 1])}^2. \end{aligned}$$

Therefore  $\|u\|_{W^{1,2}([0, 1])}^2 \geq 3/4 > 1/4$ . Part (ii) is proven.

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